

# Arvind Borde / PHY 19, Week 11: Hydrogen Again

## §11.1 The Schrödinger Equation in 3d

In three-dimensions, the  $\mathcal{S}$  eq. is

$$-\frac{\hbar^2}{2m} \nabla^2 \Psi + U(x, y, z) \Psi = i\hbar \frac{\partial \Psi}{\partial t}$$

where

$$\nabla^2 \equiv \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2}$$

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Here, too, if we assume that

$$\Psi(x, y, z, t) = \psi(x, y, z) \phi(t)$$

then we get

$$-\frac{\hbar^2}{2m} \nabla^2 \psi + U(x, y, z) \psi = E \psi$$

$$i\hbar \frac{\partial \phi}{\partial t} = E \phi$$

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## §11.2 Spherical Symmetry

A convenient coordinate to use here is the usual radial coordinate  $r$ , given by

$$r^2 = x^2 + y^2 + z^2$$

(1) Find  $\partial r / \partial x$ .

$$2r \frac{\partial r}{\partial x} = 2x$$

So

$$\frac{\partial r}{\partial x} = \frac{x}{r}$$

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The same holds for  $y$  and  $z$ .

(2) If  $f = f(r)$  is a function of  $r$ , what is  $\partial f / \partial x$ ?

$$\frac{\partial f(r)}{\partial x} = \frac{df(r)}{dr} \frac{\partial r}{\partial x} = \frac{x}{r} \frac{df(r)}{dr}$$

Assuming that  $\psi$  is a function of  $r$  alone,

$$\frac{\partial \psi}{\partial x} = \frac{x}{r} \frac{d\psi}{dr}$$

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(3) Find  $\partial^2 \psi / \partial x^2$ .

$$\begin{aligned} \frac{\partial^2 \psi}{\partial x^2} &= \frac{\partial}{\partial x} \left( \frac{x}{r} \right) \frac{d\psi}{dr} + \frac{x}{r} \frac{\partial}{\partial x} \frac{d\psi}{dr} \\ &= \left( \frac{1}{r} - \frac{x}{r^2} \cdot \frac{x}{r} \right) \frac{d\psi}{dr} + \frac{x^2}{r^2} \frac{d^2 \psi}{dr^2} \\ &= \frac{1}{r} \frac{d\psi}{dr} - \frac{x^2}{r^3} \frac{d\psi}{dr} + \frac{x^2}{r^2} \frac{d^2 \psi}{dr^2} \end{aligned}$$

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(4) Find  $\nabla^2 \psi = \frac{\partial^2 \psi}{\partial x^2} + \frac{\partial^2 \psi}{\partial y^2} + \frac{\partial^2 \psi}{\partial z^2}$ .

$$\begin{aligned} \nabla^2 \psi &= \frac{3}{r} \frac{d\psi}{dr} \\ &\quad - \frac{x^2 + y^2 + z^2}{r^3} \frac{d\psi}{dr} \\ &\quad + \frac{x^2 + y^2 + z^2}{r^2} \frac{d^2 \psi}{dr^2} \\ &= \frac{2}{r} \frac{d\psi}{dr} + \frac{d^2 \psi}{dr^2} \end{aligned}$$

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ADDITIONAL NOTES

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Let  $f' \equiv df/dr$ .

(5) What's  $(r\psi)'$ ?

$$\psi + r\psi'$$

(6) What's  $(r\psi)''$ ?

$$2\psi' + r\psi''$$

(7) Recognize this?

$$\text{It's } r\nabla^2\psi.$$

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Let  $\hat{\psi}(r) = r\psi(r)$ .

(8) Write the 3-d time-ind.  $S$  eq.

$$-\frac{\hbar^2}{2m}\nabla^2\psi(r) + U(x, y, z)\psi(r) = E\psi(r)$$

in terms of  $\hat{\psi}(r)$ . (Hint: Multiply the above by  $r$ .)

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$$r\left(-\frac{\hbar^2}{2m}\nabla^2\psi(r) + U(x, y, z)\psi(r)\right) = E(r\psi(r)) = E\hat{\psi}(r)$$

$$-\frac{\hbar^2}{2m}\hat{\psi}''(r) + U(r)\hat{\psi}(r) = E\hat{\psi}(r)$$

All this is true for any spherically symmetric case.

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### §11.3 The Hydrogen Atom

Since the proton is around a thousand times as massive as the electron (similar to the sun-Jupiter system), we'll consider this as a problem of an electron moving in an electrostatic potential created by a fixed proton.

The potential is only a function of the distance between the electron and the proton.

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For the Hydrogen atom,

$$U(r) = -k\frac{e^2}{r}$$

The equation for  $\hat{\psi}$  becomes

$$-\frac{\hbar^2}{2m_e}\hat{\psi}''(r) - k\frac{e^2}{r}\hat{\psi}(r) = E\hat{\psi}(r)$$

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Let's try  $\hat{\psi}(r) = re^{-\alpha r}$ .

(9) Find  $\hat{\psi}'(r)$ .

$$\hat{\psi}'(r) = e^{-\alpha r} - r\alpha e^{-\alpha r}$$

(10) Find  $\hat{\psi}''(r)$ .

$$\hat{\psi}''(r) = -2\alpha e^{-\alpha r} + r\alpha^2 e^{-\alpha r}$$

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#### ADDITIONAL NOTES

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(11) Plug this into the  $\mathcal{S}$  eq.

$$-\frac{\hbar^2}{2m_e} (-2\alpha e^{-\alpha r} + r\alpha^2 e^{-\alpha r}) - k\frac{e^2}{r} r e^{-\alpha r} = E r e^{-\alpha r}$$

For this to be a respectable solution, it must work for all  $r$ .

(12) What's a good  $r$  to try?

$$r = 0$$

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(13) So, what's  $\alpha$ ?

$$\frac{\hbar^2}{m_e} \alpha - k e^2 = 0$$

$$\alpha = \frac{k m_e e^2}{\hbar^2}$$

(14) Does this look familiar? Yesssss!!!!

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$\alpha^{-1}$  is the first Bohr radius.

The associated  $E$  is the ground state Bohr energy.

We get  $\psi(x) = C\hat{\psi}(x)/r = C e^{-r/r_1}$ ,  
where  $C$  is fixed by normalization.

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### §11.4 Remarks on Helium

We model this, again, as a fixed nucleus with two electrons buzzing around it.

The potential energy is

$$-k\frac{2e^2}{r_1} - k\frac{2e^2}{r_2} + k\frac{e^2}{r_{12}}$$

(15) Can you interpret these terms? Yes.

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The problem can be treated exactly if we ignore the electron-electron interaction.

It can be treated approximately, with that term.

The analysis gives energy levels that agree with observation.

This may also be done with more complicated atoms.

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ADDITIONAL NOTES

